

## Reply to "Comment on 'Uniaxial anisotropic flux trapping in Y-Ba-Cu-O and Bi-Sr-Ca-Cu-O single crystals'"

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By comparing data of the angular dependence of the remanent magnetization for a ceramic sample and for a single crystal of  $\text{YBa}_2\text{Cu}_3\text{O}_7$  with similar demagnetization factors, we conclude that intrinsic anisotropy contributes significantly to anisotropic flux trapping in the crystal.

In Ref. 1 we demonstrated the existence of uniaxial anisotropy in crystals of high-temperature superconductors (HTSC) by measuring the angular dependence of the remanent magnetization. In these measurements the sample is cooled in a field  $\mathbf{H}$ ; the "cooling angle"  $\phi$  is the angle between  $\mathbf{H}$  and a principal crystalline axis. (Here we refer only to the crystalline  $c$  axis.) At temperature  $T < T_c$  the field is turned off and the remanent magnetization  $M_{\text{rem}}$  is recorded while the sample is rotated relative to  $\mathbf{H}$ . The axis of rotation is parallel to  $c$  and  $\theta$  is the angle between  $\mathbf{H}$  and  $c$ . The angular dependence data

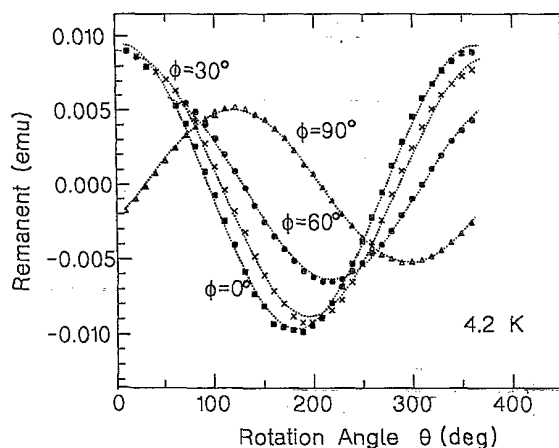


FIG. 1. Angular dependence of the remanent magnetization for the indicated cooling angles  $\phi$ , at 4.2 K and a field of 500 Oe, for a flat ceramic Y-Ba-Cu-O sample.  $\phi=0^\circ$  for  $\mathbf{H}$  parallel to the thin dimension.

( $M_{\text{rem}}$  vs  $\theta$ ), for various  $\phi$ , are exhibited in Fig. 1 of Ref. 1. The dominant features of this figure are the pronounced minima at fixed angle  $\theta=180^\circ$ , independent of  $\phi$  (except for  $\phi$  very close  $90^\circ$ ). We thus concluded that  $c$  is an anisotropic axis. In similar measurements<sup>2</sup> on ceramic HTSC, we had found these minima to be always located  $180^\circ$  relative to the cooling angle, thus  $\mathbf{H}$  defines an anisotropic direction.

Kolešník *et al.*<sup>3</sup> argue that strong demagnetization corrections of the applied field, due to the shape of the sample, induce the observed angular dependence in crystals. Their Fig. 2 demonstrates that, indeed, geometry plays an important role in the measured anisotropy. In order to estimate the relative importance of the geometry, we repeat the angular dependence experiment of Ref. 1 for a ceramic Y-Ba-Cu-O sample ( $2.1 \times 1.7 \times 0.15 \text{ mm}^3$ ) with demagnetization factors very similar to those of our Y-Ba-Cu-O crystal of Ref. 1. Here, the "cooling angle"  $\phi=0^\circ$  for  $\mathbf{H}$  parallel to the thin (0.15 mm) dimension. The results, which are shown in Fig. 1, demonstrate clearly the effect of geometry. The minima which are expected<sup>2</sup> at  $\theta=\phi+180^\circ$ , are found now at smaller  $\theta$  values, as predicted by Kolešník *et al.* However, if geometry plays the main role in determining  $M_{\text{rem}}(\theta)$  of Y-Ba-Cu-O crystals,<sup>1</sup> one would expect a similar shift in the location of the minima in these crystals. The fact that the peaks are  $\phi$  independent (Fig. 1 of Ref. 1) points to a strong intrinsic anisotropy in Y-Ba-Cu-O crystals.

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<sup>1</sup>I. Felner, U. Yaron, Y. Yeshurun, G. V. Chandrashekar, and F. Holtzberg, Phys. Rev. B **40**, 5239 (1989).

<sup>2</sup>Y. Wolfus, Y. Yeshurun, and I. Felner, Phys. Rev. B **37**, 3667

(1988).

<sup>3</sup>S. Kolešník, T. Skośkiewicz and J. Igalson, preceding Comment, Phys. Rev. B **43**, 13 679 (1991).