GRAPHITE INTERCALATION COMPOUNDS WITH ReF₆ AND TcF₆: ELECTRONIC AND MAGNETIC PROPERTIES

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ABSTRACT

Magnetic, structural and electronic properties of new graphite compounds with ReF₆ and TcF₆ are presented. Together with previous data, these provide evidence for the existence of a correlation between the layer spacing, the charge transfer and electron affinities of the transition metal hexafluorides.

INTRODUCTION

The intercalation of transition metal hexafluorides (MF₆) into graphite is of interest, because these hexafluorides comprise a unique group of isostructural and nearly isodimensional intercalants differing mainly in their electronic configurations and oxidizing properties [1-5]. The intercalation reaction may generally be written as

\[ nC + MF₆ \rightarrow Cₙ^{+x}(MF₆)ₓ(MF₆)₁₋ₓ \]  

where \( x \) denotes the degree of charge transfer. Recently [4,5], we demonstrated that in some cases \( x \) can be extracted from simultaneous ESR and susceptibility studies of the magnetic intercalate species. Particularly, for OsF₆ we confirm an earlier suggestion by Bartlett [3] that \( x = 1 \). Similar measurements on graphite compounds with MoF₆ lead to a charge transfer, \( x = 0.2 \) [5].

We report the intercalation of highly oriented pyrolytic graphite (HOPG) with the hexafluorides ReF₆ and TcF₆ for the first time.
EXPERIMENTAL

Materials

The graphite used was HOPG in the form of blocks approximately 3 x 7 x 0.2 mm. Rhenium and technetium hexafluorides were prepared by fluorination of the metal powders— the former in the presence of excess metal, the latter with excess fluorine.

Sample preparation

Intercalated samples were obtained by exposing the HOPG to the MF₆ at room temperature. Reaction rates were controlled by maintaining the MF₆ reservoirs at suitably low temperatures. The intercalation of ReF₆ works best by immersing the HOPG in liquid ReF₆ (m.p. 18.5°C). Even then stage I is reached only after about three weeks. Appropriate precautions were taken with the radioactive TcF₆ (T½ = 2 x 10 years). Stoichiometries were determined by weight uptake. Stages were characterized by (00l) diffractograms using Cu Kα radiation through Mylar covered samples. These were handled in an inert atmosphere Drilab.

Magnetic measurements

ESR measurements were carried out in the X band over the temperature range between 2K and 300K. Susceptibility measurements were performed on a superconducting quantum interference device (SQUID) down to 4K.

Electrical conductivity measurements

Basal plane electrical conductivities were measured by the contactless inductive rf technique [3]. The specific conductivity normalized per graphite plane was monitored continuously by a motor driven ferrite core and the sample in/out difference, ΔV, recorded as a function of time.

RESULTS AND DISCUSSION

Contrary to earlier reports, ReF₆ intercalates into HOPG, albeit extremely slowly. In this sense it resembles the intercalation of MoF₆ [5]. During the course of the intercalation the electrical conductivity increases smoothly reaching an asymptotic value of σ/σᵢ = 6.5 at stage I (σᵢ is the basal plane conductivity of pure HOPG). Stage I of HOPG/ReF₆ has a filled layer spacing of d = 8.35 Å; i.e. close to that of CsMoF₆ which ranges between 8.35 - 8.45 Å. Magnetic susceptibility measurements on C₁₂ReF₆ (stage I) indicate a Curie law with an effective magnetic moment of μ = 1.25 BM. This moment is between that of neutral ReF₆ (μ = 0.25 BM) [6] and that of ReF₆ in the alkali hexafluororhenates (μ = 1.5 to 2.1 BM) [7]. No ESR signal associated with ReF₆ (5d² configuration) could be observed down to T = 2K. This is not surprising as to the best of our knowledge no resonance associated with a localized spin d²...
configuration has ever been observed in any metallic environment. The lower stage HOPG/ReF₆ compounds are unstable. In the absence of excess ReF₆, they deintercalate slowly reaching a stable stage IV compound (d = 18.33Å) after about 16 hours.

The intercalation of technetium hexafluoride (TcF₆) into HOPG resembles that of OsF₆. When HOPG is exposed to TcF₆ vapor (~10 torr), the conductivity rises rapidly, reaching a maximum (σ/σ₀ ~ 7) at approximately stage II (d = 11.46 Å). Upon further exposure, stage I is reached (d = 8.12 Å) for which the conductivity ratio σ/σ₀ ~ 1. Changes in the c-axis thickness, t, confirm the values calculated from the X-ray diffraction: i.e. for stage I, the thickness ratio with respect to that of HOPG:t/t₀ = 2.4 (the calculated value based on X-rays is 2.42); for stage II, t/t₀ = 1.77 (calculated 1.71). The magnetic susceptibility of C₁₀TcF₆ (stage I) follows Curie behavior down to 4K with an effective magnetic moment μ ≈ 0.7 BM. This falls between the moments of neutral TcF₆ (μ = 0.45 BM) [6] and those of TcF⁺ in the alkali hexafluorotechnetates (μ = 2.25 to 2.51 BM) [8]. ESR studies of stage I and II HOPG/TcF₆ samples indicate clear ten-line hyperfine resolved spectra associated with ⁹⁹Tc nuclei (I = 9/2, 100% abundance) in the liquid helium temperature range. These spectra can be interpreted using a spin Hamiltonian with S = 1/2, g = 2 and isotropic hyperfine interaction. The large hyperfine interaction ⁹⁹A = 330±15 G is in agreement with previous estimates for ⁹⁹Tc in a d¹ configuration [9].

For the case of HOPG/ReF₆, the magnetic properties are consistent with the oxidizing power of ReF₆. Assuming that the effective magnetic moment of the intercalated ReF⁻⁶ is similar to that in KReF₆ (i.e. 2.1 BM) and that the magnetic properties of C₉ReF₆ are due to the presence of both neutral ReF₆ and charged ReF⁻⁶, we calculate the degree of charge transfer to be x ~ 0.35. This incomplete charge transfer is correlated with the light blue color of the final product and the very long intercalation process with the relatively high final σ as has been observed with MoF₆. The situation is more complicated in the case of TcF₆. The known chemistry of TcF₆ [10] indicates that it falls between OsF₆ and IrF₆ in oxidizing power, but probably nearer the former (see Figs 1 and 2 for the electron affinities of the MF₆ compounds). Thus, the degree of charge transfer is expected to be about one charge per TcF₆. Indeed, the dark blue color of the TcF₆/HOPG, the relatively fast intercalation, and low final σ, strongly support significant charge transfer as expected for the TcF₆ oxidizing power. However, the magnetic properties seem to be in conflict with such a conclusion. The relatively small (average) magnetic moment of HOPG/TcF₆ seems to indicate that most of the intercalant species are in the form of TcF₆ molecules. This is consistent with the ESR result, as the ESR spectra can be interpreted in terms of a 5d¹ (g=2, S=½) electronic configuration corresponding to neutral TcF₆. The magnetic properties of HOPG/TcF₆ are not understood at the present time. However,
The most striking feature of our experimental results is the existence of a correlation between the layer spacing and the charge transfer for the hexafluoride series (Fig. 1). The charge transfer was determined from the magnetic properties in the case of MoF$_6$, ReF$_6$, OsF$_6$, and PtF$_6$. For TcF$_6$ and IrF$_6$ the oxidizing properties suggest $x = 1$, and this is indicated in the figure. The plot shows that the layer spacing decreases smoothly with increasing charge transfer (or electron affinity). This can be understood by a simple classical model involving a balance between coulombic attractive and elastic repulsive interactions. The results in Fig. 1 provide evidence that the degree of charge transfer is directly reflected by the lattice spacing in the hexafluoride series. The relationship between oxidizing powers and electron affinity (E.A.) of the transition metal hexafluorides has already been pointed out [11]. For the 5d series, the E.A. were estimated by interpolation between the measured value for WF$_6$ (104 kcal-mole$^{-1}$) and one calculated for PtF$_6$ from a Born-Haber cycle for O$_2$PtF.$^-^2$.

As seen from Fig. 2, the present work puts these values on a firmer basis. It moreover indicates that the E.A. given previously for IrF$_6$ is overestimated and/
or that IrF₆ is partially reduced to IrF₆⁻² upon intercalation, and thus x > 1. The slightly shorter d-spacing for CₓIrF₆ supports this. Indeed preliminary magnetic susceptibility measurements on HOPG/IrF₆ show a slight Curie component superimposed on the large temperature independent paramagnetism reported earlier [3].

ACKNOWLEDGMENTS

We gratefully acknowledge the supply of HOPG from Dr. W.A. Moore of Union Carbide Co and financial support from the Volkswagen Foundation.

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