

# Annealing study of Fe<sub>2</sub>O<sub>3</sub> nanoparticles: Magnetic size effects and phase transformations

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Sonochemically synthesized Fe<sub>2</sub>O<sub>3</sub> nanoparticles were annealed in air or in vacuum while their magnetization was continuously recorded. Annealing in vacuum at temperatures  $T_a$  between 240 and 450 °C produced nanophases of  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub> with average particle size ranging from 4 to 14 nm, depending on  $T_a$ . Phase transformation into  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> occurred directly by annealing in air, or via an intermediate Fe<sub>3</sub>O<sub>4</sub> phase by annealing in vacuum at temperatures higher than 450 °C. Mapping the correlation between the magnetic properties and the annealing conditions, enables control of the annealing process to obtain nanocrystals of  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub>,  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>, or Fe<sub>3</sub>O<sub>4</sub> with different particle size and magnetic properties. © 2002 American Institute of Physics. [DOI: 10.1063/1.1457544]

## I. INTRODUCTION

Metal oxides are of interest to many scientific and technological disciplines.<sup>1</sup> In particular, nanostructures of these materials have attracted considerable interest as they exhibit materials properties that differ strongly from those of the bulk phases.<sup>2–6</sup> These particle size effects enable tailoring the materials to a wide range of applications, including magnetic ferrofluids, electronics and catalysis.<sup>7,8</sup> Of special interest are gamma ferric oxide ( $\gamma$ -Fe<sub>2</sub>O<sub>3</sub>) particles as they retain important position in magnetic storage media.<sup>9,10</sup>

In this article we report on results of annealing sonochemically synthesized  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub> nanocrystals in different conditions, focusing on two aspects: (1) changes in the particle size distribution, and the associated variations in magnetic properties of the material during the annealing process and (2) transformations of  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub> nanoparticles to other iron–oxide phases caused by annealing at different conditions. Our experiments show that crystallites of  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub> with different average particle size ranging from 4 to 14 nm can be obtained by controlled annealing in vacuum. We report on the continuous change in the magnetic properties of the material during this annealing process.

A well known characteristic of the iron–oxide system is the variety of possible interconversions between the different phases.<sup>1,11</sup> This has been widely investigated for large particles. Our experiments show unique transformation properties for nanosize particles of  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub> caused by annealing at different conditions (ambient, temperature and time). The identification of the various phases has been based on their color, magnetic properties and Mössbauer spectra. Mapping the correlation between the magnetic properties of the an-

nealing products and the annealing conditions, enables control of the annealing process to obtain nanocrystals of  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>,  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub> and Fe<sub>3</sub>O<sub>4</sub> with different particle size and magnetic properties.

## II. EXPERIMENT

High purity Fe<sub>2</sub>O<sub>3</sub> nanoparticles were prepared by a sonochemical method.<sup>12,13</sup> Namely, a solution of Fe(CO)<sub>5</sub> in decalin was ultrasonically irradiated under 1.5 atm of air at 0 °C for 3 h using a Sonics and Materials VC-1500 ultrasound processor. The volume of the sonicated solution was 500 ml. At the end of the sonification the solid product was separated by centrifugation, and washed thoroughly with dry pentane. The typical amount of the as-prepared product was 2.5 g. The product had a brown color typical of  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub>. As prepared and annealed samples were characterized by x-ray diffraction (XRD), transmission electron microscopy (TEM), and Mössbauer spectroscopy (MS). The <sup>57</sup>Fe MS were measured with a 50 mCi<sup>57</sup>Co:Rh source, using a constant acceleration mode. The isomer shift values are relative to Fe metals at room temperature. Magnetic properties of the samples were recorded in a vibrating sample magnetometer (Oxford 3001). Magnetic hysteresis loops were measured by applying a maximum magnetic field of 1.6 T. The saturation magnetization was obtained from magnetization versus 1/H curves, by extrapolating to 1/H=0. TEM images were utilized to measure the particle size distribution and the average particle size, defined as  $D = \langle dw \rangle / \langle w \rangle = \langle dV \rangle / \langle V \rangle = \langle d^4 \rangle / \langle d^3 \rangle$ , where  $d$ ,  $w$  and  $V$  are particle size, weight and volume, respectively. Differential scanning calorimetry (DSC) (Mettler DSC 25) at 10 °C/min was used to characterize the crystallization behavior of the material. Annealing experiments were performed in either air or vacuum maintained by a rotation pump.

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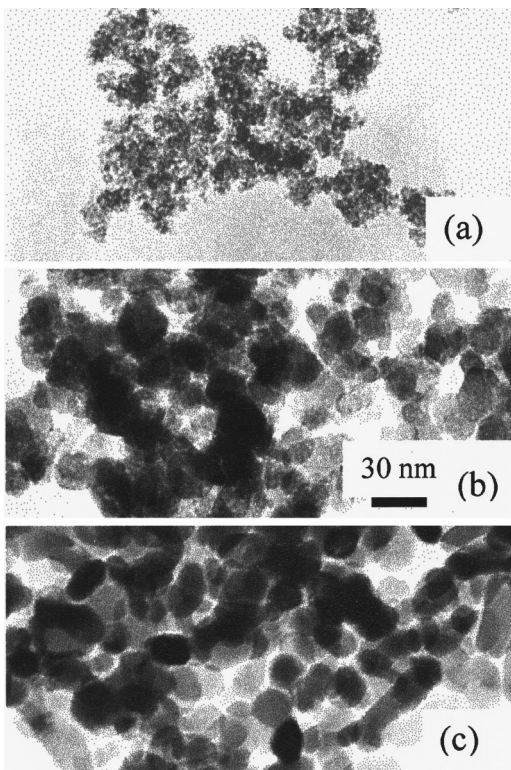


FIG. 1. RT TEM pictures of: (a) as prepared  $\text{Fe}_2\text{O}_3$ , (b)  $\gamma\text{-Fe}_2\text{O}_3$  obtained after annealing in vacuum at  $400^\circ\text{C}$  for 6 h, and (c)  $\alpha\text{-Fe}_2\text{O}_3$  obtained after sweeping temperature to  $570^\circ\text{C}$  in air.

### III. CHARACTERIZATION OF AS-PREPARED SAMPLES

Figure 1(a) exhibits a TEM image of the as-prepared material, indicating an average particle size  $D \sim 3$  nm with quite narrow particle size distribution [see Fig. 2(a)]. The smeared XRD peak shown in Fig. 3(a) is consistent with this value of  $D$ . Figure 3(b) shows the magnetization curve measured at room temperature (RT). As expected, it exhibits a superparamagnetic behavior, characteristic of small magnetic particles.<sup>2-4</sup> The magnetization curve is practically reversible with negligible coercive field  $H_c = 0.5$  Oe and remanent magnetization  $M_R = 5 \times 10^{-4}$  emu/g. DSC, Fig. 3(c), shows a smeared exothermic peak around  $258^\circ\text{C}$ , with full width at half maximum from  $160$  to  $310^\circ\text{C}$ , indicating an increase in the average size of the crystallites.

The RT Mössbauer spectrum, Fig. 4(a), consists of a broad doublet indicating the absence of a long range magnetic ordering. Computer analysis reveals the presence of two quadrupole doublets with the same isomer shift ( $IS$ ) =  $0.46$  mm/s and linewidth of  $W = 0.44(1)$  mm/s. The relative ratio of 4:1 corresponds to inequivalent Fe sites in the material. The quadrupole splitting (QS) are:  $QS = eqQ/2 = 1.19(1)$  and  $0.69(1)$  for the major and minor doublets, respectively. These values are typical of  $\text{Fe}^{+3}$  in the high spin state and are consistent with values obtained in Mössbauer spectra of  $\text{Fe}_2\text{O}_3$  nanoparticles.<sup>12,13</sup>

### IV. PARTICLE SIZE EFFECTS

As a result of annealing at elevated temperatures (above  $\sim 240^\circ\text{C}$ ) the average particle size always increases and the

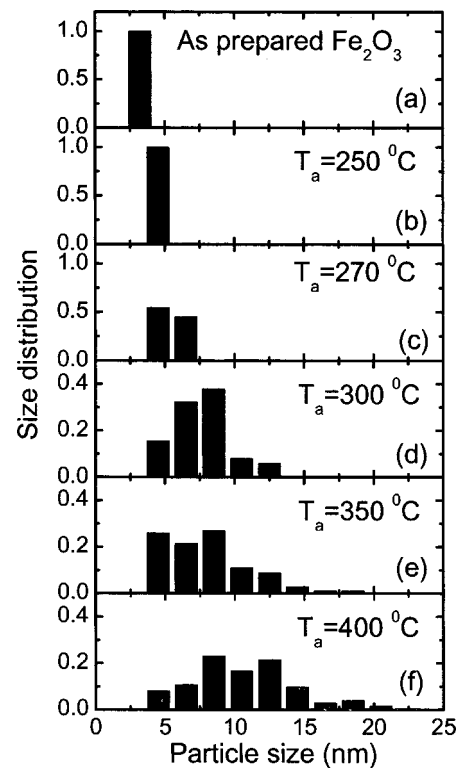


FIG. 2. Size distributions of  $\text{Fe}_2\text{O}_3$  nanoparticles: (a) as-prepared sample, (b)–(f) after annealing at  $T_a = 250, 270, 300, 350,$  and  $400^\circ\text{C}$ .

particle size distribution widens. These changes may also be accompanied by a transformation into a different iron-oxide phase, such as  $\alpha\text{-Fe}_2\text{O}_3$  or  $\text{Fe}_3\text{O}_4$  as described in Sec. V. In this section we report on annealing experiments in which the end product is always a  $\gamma\text{-Fe}_2\text{O}_3$  nanophase with larger average particle size. This result has been obtained in vacuum annealing at temperatures from  $240$  to  $450^\circ\text{C}$  for time periods from 2 to 6 h. During the annealing process, the magnetization of the samples was continuously recorded.

Figure 5(a) shows the isothermal normalized magnetization  $M/M_0$  versus annealing time  $t_a$  at annealing temperatures  $T_a$  between  $240$  and  $450^\circ\text{C}$ , in the presence of a magnetic field of  $500$  Oe. ( $M_0 = 0.37$  emu/g is the RT magnetization of the as prepared sample). Evidently, the magnetization increases with the annealing time  $t_a$  until a stable saturation value  $m_s$  is reached. The annealing products were characterized at room temperature by magnetic hysteresis loop measurements, Mössbauer spectroscopy, and TEM.

The magnetization and the coercive field increase monotonically with annealing temperature as shown in Figs. 6(a) and 6(b). We ascribe this behavior to the increase in the average crystallite size of the  $\gamma\text{-Fe}_2\text{O}_3$  nano-phase caused by the annealing. This claim has been confirmed by RT Mössbauer spectra of the final products and TEM measurements as described below.

Figures 4(b)–4(f) show the RT  $^{57}\text{Fe}$  Mössbauer spectra measured after annealing at  $250 \leq T_a \leq 400^\circ\text{C}$  in vacuum. The main effect to be seen is the evolution of a doublet into a sextet through a series of complicated spectra composed of doublets and sextets. The doublet represents the fraction of small particles in which long range magnetic ordering is ab-

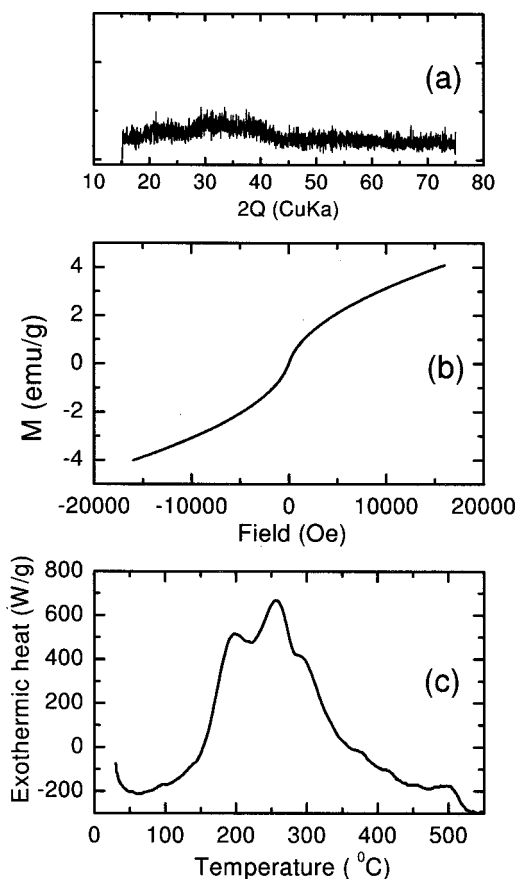


FIG. 3. (a) XRD pattern, (b) magnetization loop, and (c) DSC curve for as prepared  $\text{Fe}_2\text{O}_3$ .

sent. The fraction of the sextet increases with the annealing temperature, indicating that more particles become magnetically ordered. A detailed analysis yields the sextet fraction of 37%, 77%, 93%, and 95% for  $T_a=250$ , 270, 300, and 350 °C, respectively. TEM measurements indicate corresponding average particle size of about 4.5, 6, 9, and 12 nm, respectively. The data indicate that at 400 °C the fraction of the magnetically ordered material is close to 100%. This is consistent with the magnetic data described above. The Mössbauer spectrum measured after annealing at 400 °C (yielding 14 nm average particle size) consists of a pure sextet [see Fig. 4(f)]. Analysis of the sextet indicates a distribution of magnetic hyperfine fields due to the particle size distribution. The hyperfine parameters for this sextet are: a mean magnetic hyperfine field  $H_{\text{eff}}=52.0(5)T$ , linewidth  $W=0.40$  mm/s, and  $IS=0.48$  mm/s. The sextet does not show a quadrupole effect. These parameters are consistent with the hyperfine values reported for bulk  $\gamma\text{-Fe}_2\text{O}_3$  crystals.<sup>1,14</sup> We note that since the hyperfine field parameters for both  $\alpha\text{-Fe}_2\text{O}_3$  and  $\gamma\text{-Fe}_2\text{O}_3$  are very close to each other, and the measured lines are quite broad, the Mössbauer spectrum is not sensitive enough to conclusively determine the magnetic phase obtained. However, the large magnetization value measured for these samples clearly indicates a  $\gamma\text{-Fe}_2\text{O}_3$  phase.

The magnetization and Mössbauer measurements described above were complemented by particle size distribu-

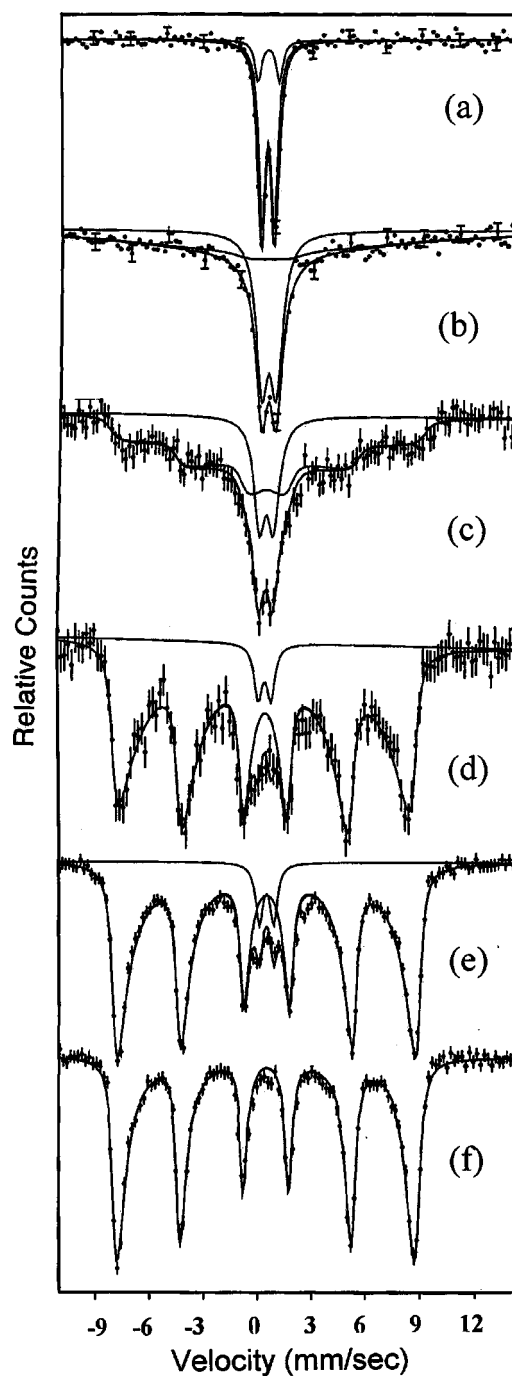


FIG. 4. RT MS of as prepared  $\text{Fe}_2\text{O}_3$  (a), and its annealing products  $\gamma\text{-Fe}_2\text{O}_3$  at various average size 4.5, 6, 9, 12, and 14 nm obtained at  $T_a=250$ , 270, 300, 350, and 400 °C, from (b) to (f), respectively.

tion measurements of the same samples after the completion of the annealing process. A typical TEM image of a sample annealed at  $T_a=400$  °C is shown in Fig. 1(b). The TEM images reflect roughly spherical nanoparticles with different size distribution depicted in Figs. 2(b)–2(f). Figure 6(c) shows the average particle size versus annealing temperature. As the annealing temperature increases from 250 to 400 °C the average particle size increases from 4.5 to 14 nm. The increase of the average particle size is accompanied by broadening of the particle size distribution as shown in Figs. 2(b)–2(f).

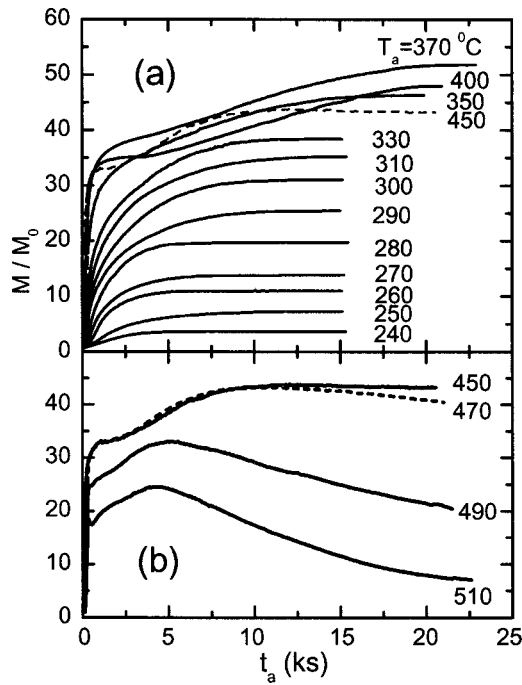


FIG. 5. Isothermal normalized magnetization vs annealing time at: (a)  $240\text{ }^\circ\text{C} \leq T_a \leq 450\text{ }^\circ\text{C}$ , and (b) for  $T_a > 450\text{ }^\circ\text{C}$ .

We conclude that the results of the magnetization measurements after the annealing process described in Fig. 5(a) can be fully ascribed to particle size effects. To summarize these results, in Fig. 7 we present the RT saturation magnetization  $M_s$  and coercive field  $H_c$  versus the average particle size  $D$ . Evidently,  $M_s$  and  $H_c$  drop sharply as the average particle size decreases below  $\sim 12\text{ nm}$ .<sup>15-17</sup> A fit of the ex-

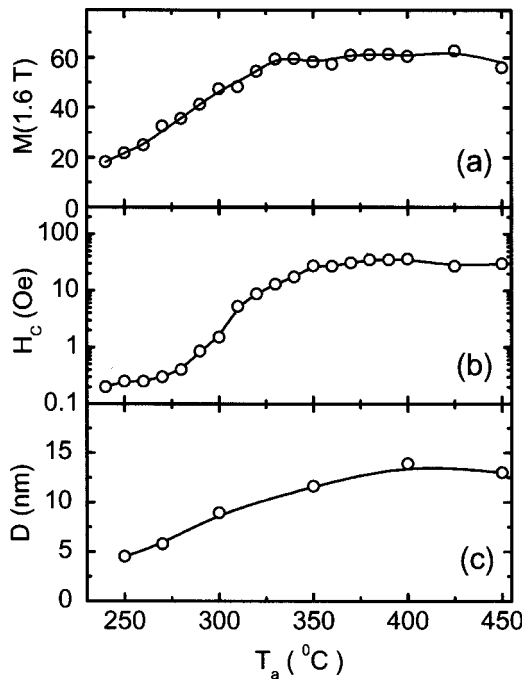


FIG. 6. (a) RT magnetization, (b) coercive field, and (c) average size of  $\gamma\text{-Fe}_2\text{O}_3$  particles obtained after annealing at various temperatures.

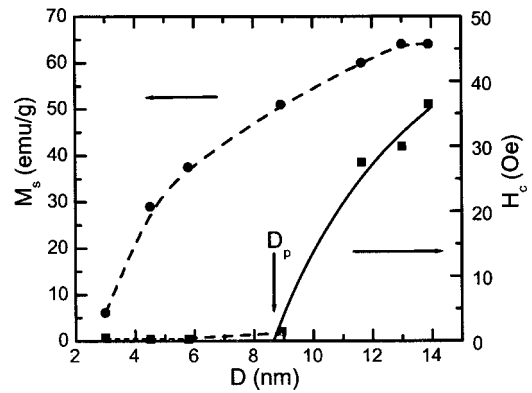


FIG. 7. RT saturation magnetization  $M_s$  and coercive field  $H_c$  vs average particle size  $D$ . Dashed lines are a guide to the eye. Solid line is a fit to  $H_c \propto [1 - (D_p/D)^{3/2}]$ .

perimental data for  $H_c$  to the relation  $H_c \propto [1 - (D_p/D)^{3/2}]$ , defines the critical size  $D_p \approx 8.6\text{ nm}$  below which the particles become superparamagnetic.<sup>2</sup>

V. PHASE TRANSFORMATIONS

As mentioned in the previous section, the growth of the  $\gamma\text{-Fe}_2\text{O}_3$  nanocrystals may be accompanied by transformations to other iron-oxide phases.<sup>1,18</sup> Phase transformation to  $\text{Fe}_3\text{O}_4$  occurring during the initial stages of the growth process was noticed in measurements of the magnetization of as-prepared samples versus temperature in vacuum. Figure 8 shows the magnetization  $M$  normalized to the RT magnetization  $M_0 = 2.1\text{ emu/g}$ , measured with an applied magnetic field of  $5\text{ kOe}$  in vacuum. Temperature was increased at a rate of  $10\text{ }^\circ\text{C/min}$  from  $27$  to  $640\text{ }^\circ\text{C}$ , and then decreased at a rate dictated by the natural cooling of the furnace. Upon heating, the normalized magnetization  $M(T)/M_0$  exhibits a non-monotonic behavior: initially it decreases with temperature up to approximately  $200\text{ }^\circ\text{C}$ , then it exhibits a sharp increase up to  $\sim 400\text{ }^\circ\text{C}$  followed by a decrease as the Curie temperature  $T_c \approx 570\text{ }^\circ\text{C}$  is approached. This nonmonotonic behavior can be explained as resulting from two competing processes: the normal decrease of magnetization with temperature as the Curie point is approached, and the increase in the

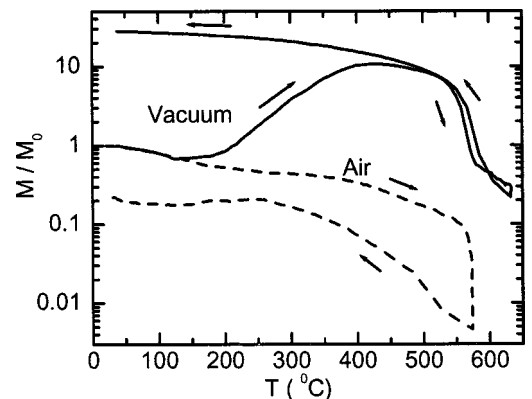


FIG. 8. Normalized magnetization vs temperature for sample in vacuum (solid line) and in air (dashed).



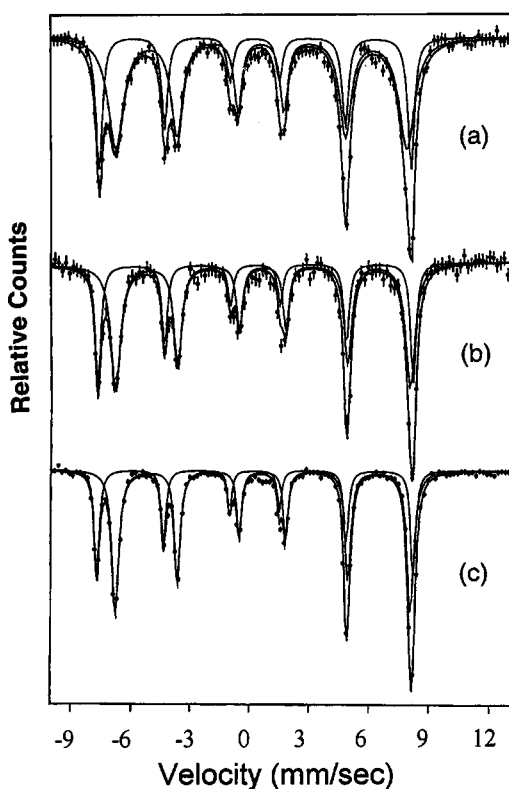


FIG. 9. RT MS of  $\text{Fe}_3\text{O}_4$  obtained by annealing as prepared  $\text{Fe}_2\text{O}_3$  at: (a)  $T_a=500^\circ\text{C}$  for 8 ks, (b)  $T_a=600^\circ\text{C}$  for 12 ks; and (c) commercial  $\text{Fe}_3\text{O}_4$ .

average particle size and magnetic ordering, as discussed in the previous section. Indeed, upon cooling the magnetization exhibits the normal behavior of the magnetically ordered phase, as no particle size changes take place. However, we noticed that during this temperature cycle a phase transformation has also occurred. This was indicated by a change of the sample color from brown to black, and further confirmed by Mössbauer spectroscopy which clearly identified the final phase as  $\text{Fe}_3\text{O}_4$ . MS of such samples are shown in Figs. 9(a) and 9(b) together with the MS of commercial  $\text{Fe}_3\text{O}_4$  [Fig. 9(c)]. The spectra shown in Figs. 9(a) and 9(b) were measured at room temperature after annealing at  $T_a=500^\circ\text{C}$  for 2 h and  $T_a=600^\circ\text{C}$  for 3 h, respectively. Comparing these spectra with the spectrum of commercial  $\text{Fe}_3\text{O}_4$  [Fig. 9(c)], one identifies similar features. Basically, these spectra consist of two superimposed sextets corresponding to two different sites of  $\text{Fe}^{+3}$  and  $\text{Fe}^{+2}$  ions. Analysis of the spectrum shown in Fig. 9(a) yields  $H_{\text{eff}}=493$  kOe,  $\text{IS}=0.70(1)$  mm/s, and  $H_{\text{eff}}=461$  kOe,  $\text{IS}=0.17(1)$  mm/s for the two sextets, respectively. These results are in agreement with the hyperfine parameters of commercial  $\text{Fe}_3\text{O}_4$ .<sup>1,19</sup> The  $\text{Fe}_3\text{O}_4$  phases corresponding to Figs. 8(a) and 8(b) exhibited a RT saturation magnetization of 87 and 93 emu/g, coercive field of 70 and 110 Oe, and average particle size of about 20 and 25 nm, respectively. These magnetization values are higher than those of  $\gamma\text{-Fe}_2\text{O}_3$  crystallites of the same size.

A different phase transformation was noticed when the samples were heated in air. The dashed line in Fig. 8 show results of magnetization measurements in air on heating and natural cooling at the same rates as described above. Upon

heating, the normalized magnetization decreases monotonically with temperature. On cooling, the magnetization follows a different curve, much lower than the curve obtained during heating. The RT product exhibits a change of color from brown to red-brown, and low normalized magnetization of  $\sim 0.2$ , i.e., much smaller than that of the as-prepared material. The TEM image of this product, [Fig. 1(c)] shows a change in the morphology of the particles from spherical to elongated shape with average length of 30 nm and a mean diameter of 18 nm. The decrease in the magnetization despite the increase in the particle size clearly indicates a phase transformation, most likely from  $\gamma\text{-Fe}_2\text{O}_3$  to  $\alpha\text{-Fe}_2\text{O}_3$ .

A phase transformation to  $\alpha\text{-Fe}_2\text{O}_3$  was also noticed by isothermal magnetization in long time annealing experiments in vacuum at temperatures above  $450^\circ\text{C}$ . The results of these experiments are shown in Fig. 5(b). For  $T_a>470^\circ\text{C}$  the magnetization decreases after  $t_a>3$  h. The final product is a red-brown colored  $\alpha\text{-Fe}_2\text{O}_3$ , with relatively small magnetization. Transformation from  $\gamma\text{-Fe}_2\text{O}_3$  to  $\alpha\text{-Fe}_2\text{O}_3$  above  $450^\circ\text{C}$  was previously noticed in different types of measurements.<sup>14,18,20</sup>

## VI. CONCLUSIONS

Nanophases of  $\gamma\text{-Fe}_2\text{O}_3$  with various average particle size can be obtained by annealing of sonochemically synthesized  $\text{Fe}_2\text{O}_3$  nanoparticles in vacuum at temperatures  $240 \leq T_a \leq 450^\circ\text{C}$  for time periods  $2\text{h} \leq t_a \leq 6\text{h}$ . The growth of the  $\gamma\text{-Fe}_2\text{O}_3$  nanocrystallites in vacuum annealing proceeds via the formation of an intermediate  $\text{Fe}_3\text{O}_4$  phase.  $\gamma\text{-Fe}_2\text{O}_3$  crystallites with average size less than 12 nm show magnetic characteristics that strongly differ from those of the bulk  $\gamma\text{-Fe}_2\text{O}_3$  phase. Annealing of  $\gamma\text{-Fe}_2\text{O}_3$  nanoparticles in air results in a direct transformation into the  $\alpha\text{-Fe}_2\text{O}_3$  nanophase. An indirect transformation into the  $\alpha\text{-Fe}_2\text{O}_3$  nanophase via intermediate  $\text{Fe}_3\text{O}_4$  and  $\gamma\text{-Fe}_2\text{O}_3$  phases may also be obtained in vacuum annealing at temperatures higher than  $450^\circ\text{C}$ .

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